

Appendix C

Derivation of the Bogoliubov-de Gennes equations

Here we present a derivation of the Bogoliubov-de Gennes equations for the dilute Bose gas, describing the spectrum of small-scale excitations on the background state.

If $\phi_0(\mathbf{r})$ is the time-independent background state, then weak perturbations of the wavefunction can be expressed with the trial wavefunction,

$$\psi(\mathbf{r}, t) = e^{-i\mu t/\hbar} \left[\phi_0(\mathbf{r}) + u(\mathbf{r})e^{-i\omega t} + v^* e^{i\omega t} \right], \quad (\text{C.1})$$

where μ is the chemical potential of the system. The perturbation is time-dependent with characteristic frequency ω . The amplitude functions $u(\mathbf{r})$ and $v(\mathbf{r})$ are complex. For simplicity, we will drop the time and space notation in the parameters.

The time-dependent Gross-Pitaevskii equation is,

$$i\hbar \frac{\partial \psi}{\partial t} = \left[H_0 + g|\psi|^2 \right] \psi, \quad (\text{C.2})$$

where H_0 is the single particle hamiltonian $H_0 = -(\hbar^2/2m)\nabla^2 + V_{\text{ext}}$, with V_{ext} being the external potential acting on the system, and g is the nonlinear coefficient given by equation (2.7).

We proceed by evaluating the Gross-Pitaevskii equation for the trial wavefunction of equation (C.1) to first order in u and v . The left-hand side of the

Gross-Pitaevskii equation becomes,

$$i\hbar \frac{\partial \psi}{\partial t} = e^{-i\mu t/\hbar} \left[\mu \phi_0 + (\mu + \hbar\omega) u e^{-i\omega t} + (\mu - \hbar\omega) v^* e^{i\omega t} \right]. \quad (\text{C.3})$$

We now evaluate the right-hand side of the Gross-Pitaevskii equation. First we consider the term $\psi^* \psi \psi$. Keeping only the terms which are linear in u and v gives,

$$\begin{aligned} \psi^* \psi \psi &= e^{-i\mu t/\hbar} (\phi_0^* + u^* e^{i\omega t} + v e^{-i\omega t}) (\phi_0 + u e^{-i\omega t} + v^* e^{i\omega t})^2 \\ &= e^{-i\mu t/\hbar} \left[|\phi_0|^2 (\phi_0 + 2u e^{-i\omega t} + 2v^* e^{i\omega t}) + \phi_0^2 (u^* e^{i\omega t} + v e^{-i\omega t}) \right]. \end{aligned} \quad (\text{C.4})$$

Substituting all terms into the Gross-Pitaevskii equation and equating terms in $e^{-i\mu t/\hbar}$ leads to,

$$\mu \phi_0 = H_0 \phi_0 + g |\phi_0|^2 \phi_0. \quad (\text{C.5})$$

Equating terms in $e^{-i(\mu+\hbar\omega)t/\hbar}$ gives,

$$(\mu + \hbar\omega) u = (H_0 + 2g |\phi_0|^2) u + g \phi_0^2 v, \quad (\text{C.6})$$

and terms in $e^{-i(\mu-\hbar\omega)t/\hbar}$ gives,

$$(\mu - \hbar\omega) v^* = (H_0 + 2g |\phi_0|^2) v^* + g \phi_0^2 u^*. \quad (\text{C.7})$$

Taking the complex conjugate of equation (C.7) leads to,

$$(\mu - \hbar\omega) v = (H_0 + 2g |\phi_0|^2) v + g \phi_0^{*2} u. \quad (\text{C.8})$$

Equations (C.5), (C.6) and (C.8) can be written in the form of a matrix equation in terms of u and v ,

$$\begin{pmatrix} L & M \\ M^* & L \end{pmatrix} \begin{pmatrix} u \\ v \end{pmatrix} = \hbar\omega \begin{pmatrix} u \\ -v \end{pmatrix}, \quad (\text{C.9})$$

where,

$$L = H_0 - \mu + 2g |\phi_0|^2, \quad (\text{C.10})$$

and,

$$M = g \phi_0^2. \quad (\text{C.11})$$

This is the Bogoliubov-de Gennes equation for the dilute Bose gas. Solutions to this matrix equation can be obtained by solving the eigenvalue equation,

$$\begin{vmatrix} L - \hbar\omega & M \\ M^* & L + \hbar\omega \end{vmatrix} = 0. \quad (\text{C.12})$$

It follows that the energy of the excitations is given by,

$$\hbar\omega = \sqrt{L^2 - |M|^2}. \quad (\text{C.13})$$

In particular, for a homogeneous system with background density n , $V_{\text{ext}} = 0$, $H_0 = \hbar^2 k^2 / 2m$, and $\mu = ng$, where k is the wavevector of the excitation. The energy spectrum is then given by,

$$\hbar\omega = \sqrt{\frac{\hbar^2 k^2}{2m} \left(\frac{\hbar^2 k^2}{2m} + 2ng \right)}. \quad (\text{C.14})$$

This dispersion is particle-like ($\hbar\omega = \hbar^2 k^2 / 2m$) at high momenta, while at low momenta, the spectrum is linear ($\omega = ck$) and supports sound (phonon) waves with characteristic speed c .